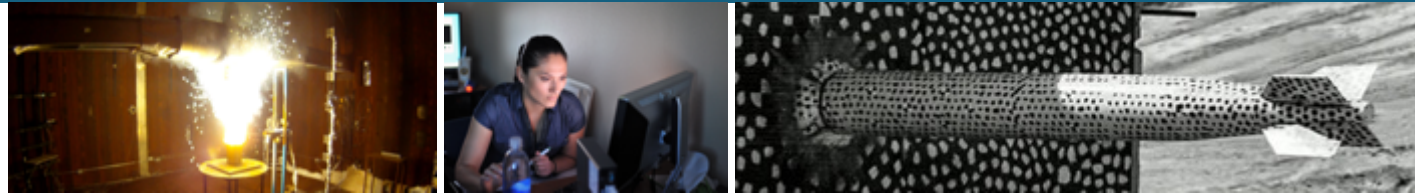


Exponential Integrators for the HOMME-NH Nonhydrostatic Dycore



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Time-stepping in HOMME-NH



The E3SM atmosphere model (EAM) targets high-resolution (3 km horizontal grid spacing) simulations. HOMME-NH uses a horizontally explicit vertically implicit (HEVI) partitioning of the nonhydrostatic (NH) equations:

$$\begin{aligned}
 \mathbf{u}_t + (\nabla_\eta \times \mathbf{u} + 2\Omega) \times \mathbf{u} + \frac{1}{2} \nabla_\eta (\mathbf{u} \cdot \mathbf{u}) + \dot{\eta} \frac{\partial \mathbf{u}}{\partial \eta} + \frac{1}{\rho} \nabla_\eta p + \mu \nabla_\eta \phi &= 0, \quad \dot{\eta} := d\eta/dt \\
 w_t + \mathbf{u} \cdot \nabla_\eta w + \dot{\eta} \frac{\partial w}{\partial \eta} + \boxed{\mathbf{g}(1 - \mu)} &= 0, \quad \mu := \frac{\partial p}{\partial \eta} / \frac{\partial \pi}{\partial \eta} \\
 \phi_t + \mathbf{u} \cdot \nabla_\eta \phi + \dot{\eta} \frac{\partial \phi}{\partial \eta} + \boxed{-\mathbf{g}w} &= 0 \\
 \Theta_t + \nabla_\eta \cdot (\Theta \mathbf{u}) + \frac{\partial}{\partial \eta} (\Theta \dot{\eta}) &= 0, \quad \Theta := \frac{\partial \pi}{\partial \eta} \theta_v \\
 \frac{\partial}{\partial t} \left(\frac{\partial \pi}{\partial \eta} \right) + \nabla_\eta \cdot \left(\frac{\partial \pi}{\partial \eta} \mathbf{u} \right) + \frac{\partial}{\partial \eta} \left(\frac{\partial \pi}{\partial \eta} \dot{\eta} \right) &= 0.
 \end{aligned}$$

- Accurate and efficient time-stepping is critical for the atmosphere to be performant.
- Semi-implicit and implicit-explicit (IMEX) methods are the “industry standard” for time-stepping in HEVI NH models.
- Exponential-type (e.g. exponential time difference (ETD) and integrating factor methods) are an attractive alternative in the HEVI context.
- Big unknown: Are nonlinear solves (IMEX) cheaper than matrix exponentials (ETD) in NH models.

ETD Runge-Kutta (ETD-RK) methods



Consider a partitioned initial value ODE $q_t = n(q) + s(q)$ where $n(q)$ represent non-stiff terms and $s(q)$ represents stiff terms.

Given an approximate solution $q_m \approx q(t_m)$ at time-step m , we linearize the stiff part $s(q)$ along q_m . Let $L_m = \partial s / \partial q(q_m)$ and let $N_m(q) = f(q) - L_m q$. Now our additive partitioning becomes:

$$q_t = f(q) = L_m q + N_m(q).$$

An ETD-RK solution is given by

$$q_{m+1} = e^{L_m \Delta t} q_m + \Delta t \sum_{i=1}^r b_i(L_m \Delta t) N_m(t_m + c_i \Delta t, g_{mi}),$$

with stage values

$$g_{mj} = e^{c_j L_m \Delta t} q_m + \Delta t \sum_{i=1}^r a_{ij}(L_m \Delta t) N_m(t_m + c_i \Delta t, g_{mi}).$$

where the method coefficients a_{ij} and b_i depend on the so-called phi-functions:

$$\varphi_0(z) = \exp(z) \quad \varphi_{k+1}(z) = z^{-1}(\varphi_k(z) - \varphi_k(0))$$

Methods are represented with the Butcher notation:

c_1	$a_{1,1}$	\dots	$a_{1,r}$
\vdots	\vdots	\ddots	\vdots

HEVI matrix exponentials in HOMME-NH



In HOMME-NH Linearizing the stiff terms results in:

$$L_m = \begin{bmatrix} 0 & & & \\ & 0 & \mathbf{g} \frac{\partial \mu}{\partial \phi} & \\ & \mathbf{g} I & 0 & \\ & & & 0 \end{bmatrix} \quad e^{\alpha L_m} = \begin{bmatrix} I & & & \\ & \exp \left(\begin{bmatrix} 0 & \mathbf{g} \alpha \frac{\partial \mu}{\partial \phi} \\ \mathbf{g} \alpha I & 0 \end{bmatrix} \right) & & \\ & & I & \\ & & & I \end{bmatrix}, \quad \alpha \neq 0$$

where $\frac{\partial \mu}{\partial \phi}$ is tridiagonal. We implement a (p, q) -Padé approximation to approximate the exponential of a matrix A as $e^A \approx [Q_{pq}(A)]^{-1} P_{pq}(A)$, where $P_{pq}(A)$ and $Q_{pq}(A)$ are defined as

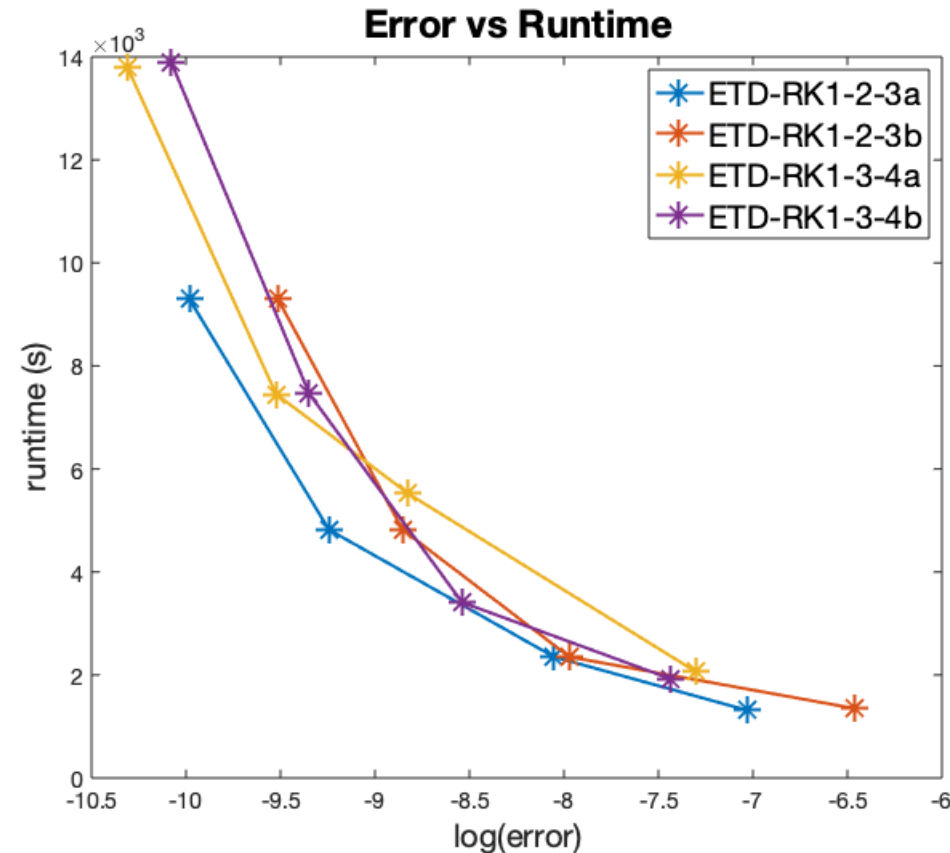
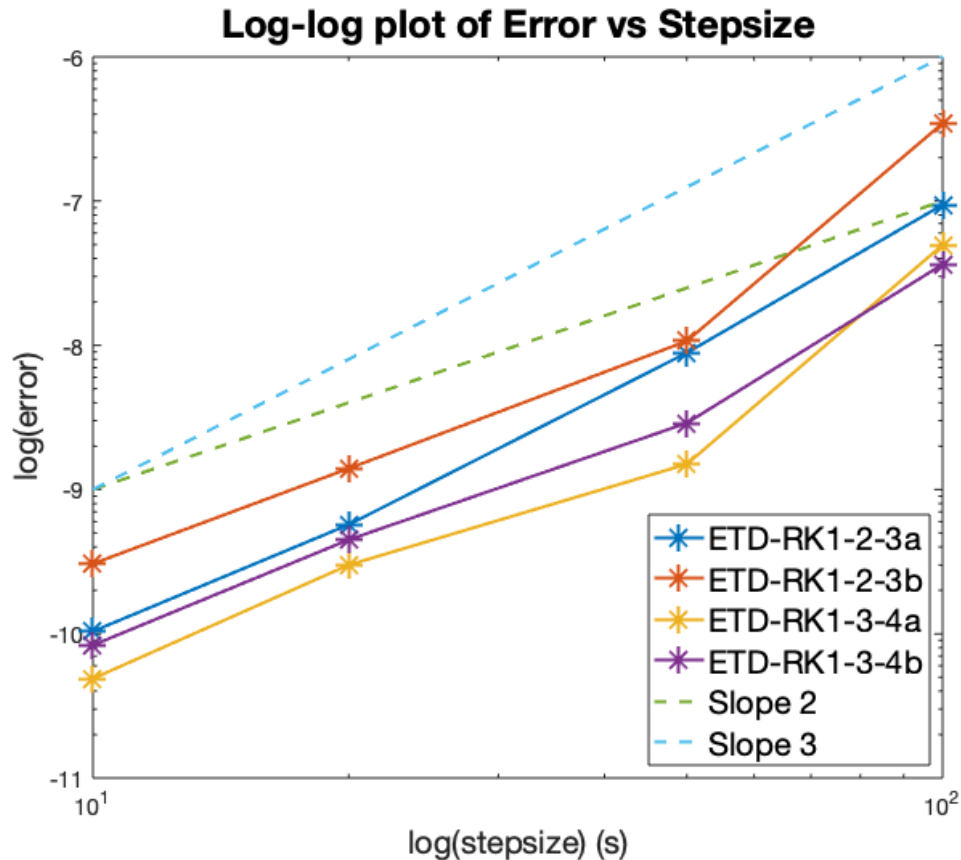
$$P_{pq}(A) = \sum_{j=0}^p \frac{(p+q-j)! p!}{(p+q)! j! (p-j)!} A^j \quad Q_{pq}(A) = \sum_{j=0}^q \frac{(p+q-j)! q!}{(p+q)! j! (q-j)!} (-A)^j$$

Factoring the polynomial $Q_{pq}(A)$ and applying static condensation to L_m means $e^{\alpha L_m}$ can be approximated with q tridiagonal solves and p tridiagonal multiplications. phi-functions φ_k with $k \geq 0$ are computed iteratively.

Taking L_m from a spin-up run of the DCMIP Test Case 3.1, the (3,3)-Padé approximation for

e^{L_m} has an absolute error of $5.25e - 13$.

Convergence and Efficiency



Convergence and efficiency study for several custom ETD-RK methods integrating the DCMIP2012 Test Case 3.1 (nonhydrostatic gravity wave). The error is the full L2-error of the solution found by comparison with a high-precision reference trajectory.

Future Work:

- Investigate parallel implementation of exponential calculations and function evaluations. Optimize performance relative to the best IMEX-RK methods.



- . Exponential Rosenbrock integrators for nonhydrostatic atmosphere models. Krause, C. and Steyer, A. To be submitted (2020)
- . Efficient IMEX methods for nonhydrostatic dynamics. Steyer,, A., Vogl, C., Taylor, M., Guba, O. **arXiv:1906.07219.**
- . Evaluation of Implicit-Explicit Additive Runge-Kutta Integrators for the HOMME-NH Dynamical Core. Vogl, C., Steyer, A., Reynolds, R., Ullrich, P., Woodward, C. **arXiv: 1904.10115.**
- . An energy consistent discretization of the nonhydrostatic equations in primitive variables. Taylor, M., Guba, O., Steyer, A., Ullrich P., Hall, D., Eldred, C. **arXiv:1908.04430.**